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S/141/60/003/005/001/026
EO 52/E 514

Determination of the Relative Fluctuations of the Electron Concentration in the Ionosphere

and radioastronomical methods. The present paper reports results of such a determination which was carried out at the Scientific Research Radiophysical Institute at Gorkiy University in the Autumn of 1959. If it is assumed that the irregularities in the electron concentration have mean linear dimensions ξ and are distributed uniformly throughout the ionospheric layer then in the case of radioastronomical observations δN is given by

$$\delta N = 0.34 \frac{\lambda^2}{\lambda_0 \sqrt{\xi z_m}} \sqrt{\ln \left(\frac{P_s}{P} + 1 \right)} \quad (1)$$

while in the pulse method it can be estimated from the inequality given by Eq. (2), where λ_0 and λ_c are the working and critical wavelength, respectively and \bar{z}_m is the effective Card 2/7

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thickness of the ionosphere P_s/P is the ratio of scattered to transmitted signal energies and I is an integral depending on the parameters of the layer The expression

$$\delta N < 0.17 \frac{\lambda_0}{\sqrt{\xi I}} \sqrt{\ln \left(\frac{P_s}{P} + 1 \right)} \quad (2)$$

holds for a signal which has been reflected only once In the case of double reflection the inequality

$$\delta N < 0.17 \frac{\lambda_0}{\sqrt{\xi I}} \sqrt{\frac{1}{2} \ln \left(\frac{P_s}{P} + 1 \right)} \quad (3)$$

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must be employed. The above formulae were derived on the assumption that

- 1) the relative fluctuations in the dielectric constant are small.
- 2) the angles of scattering θ are small ($\ll 1$).
- 3) the geometrical-optics approximation holds and
- 4) the point of observation is at a great distance from the scattering region.

It was also assumed that the dependence of the concentration on altitude is given by

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$$N = \begin{cases} N_o \left[1 - \left(\frac{z - z_o}{z_m} \right)^2 \right] & (z < z_o) \\ N_o \exp[-u(z - z_o)] & (z \geq z_o) \end{cases} \quad (4)$$

where z_o is the altitude of the maximum

N_o is the electron concentration at $z = z_o$.

Under these conditions the effective thickness of the
ionosphere is given by

$$z_m = z_o + 1/u$$

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In the calculations z_m was assumed to be equal to 400 km.

The dimensions of the irregularities ξ were estimated from the formula $\xi = VT$ where T is the mean period of fluctuations and V is the velocity of motion of the irregularities. In calculating ξ it was assumed that these irregularities move with a mean velocity of 100 m/sec. The integral I was calculated in Ref. 3. It was found that the pulse method gives $\delta N < 4 \times 10^{-3}$ and $\delta N < 10^{-2}$ for the F and E layers, respectively. Fluctuations in the emission of discrete sources (Cassiopea A and Cygnus A) gave the value of $\delta N \sim 3 \times 10^{-3}$. Acknowledgments are expressed to V.L. Ginzburg and G.G. Getmantsev for their interest and valuable advice.

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Determination of the Relative Fluctuations of the Electron
Concentration in the Ionosphere

There are 5 references 4 Soviet and 1 English

ASSOCIATION: Nauchno-issledovatel'skiy radiofizicheskiy
institut pri Gor'kovskom universitete
(Scientific Research Radiophysical Institute
of Gor'kiy University)

SUBMITTED: June 6, 1960

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9.7120 (also 1041, 1046)
9.9400

S/141/60/003/006/004/025
EO32/E114

AUTHORS: Mityakova, E.Ye., Mityakov, N.A., and Rapoport, V.O.

TITLE: On the Measurement of the Electron Concentration in the Ionosphere and in Interplanetary Space

PERIODICAL: Izvestiya vysshikh uchebnykh zavedeniy, Radiofizika, 1960, Vol.3, No.6, pp. 949-956

TEXT: A brief review is given of the available methods for the determination of the electron concentration in the ionosphere with the aid of artificial earth satellites. Using the quasi-longitudinal approximation, an expression is obtained for the phase and group paths for a signal emitted from an artificial earth satellite towards a spherical earth. It is shown using the results of Al'pert et al (Ref.11) that the phase path length is given by

$$n_{1,2}^2 = 1 - \frac{4\pi e^2 N}{m \omega (\omega \pm \omega_L)} = 1 - \frac{2aN}{\omega (\omega \pm \omega_L)} \quad (1)$$

and

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$$L_{\phi 1,2} = r_0 - \frac{a}{\omega \cos \gamma} \left[\int_0^z \frac{N}{\omega \pm \omega_L} dz - \operatorname{tg}^2 \gamma \int_0^z \frac{Nz}{R_0 (\omega \pm \omega_L)} dz \right] \quad (5)$$

and the group path length is given by

$$L_{rp1,2} = \int_A^B \frac{\partial (n_{1,2} \omega)}{\partial \omega} d\omega_{1,2} \quad (6)$$

and

$$L_{rp1,2} = r_0 + \frac{a}{\omega \cos \gamma} \left[\int_0^z \frac{N}{\omega \pm \omega_L} dz - \operatorname{tg}^2 \gamma \int_0^z \frac{Nz}{R_0 (\omega \pm \omega_L)} dz \right] \quad (7)$$

In these expressions $\omega_L = (eH_0/mc) \cos \gamma$, γ is the angle between the earth's magnetic field and the wave normal, suffix 2 and the "minus" sign refer to the ordinary wave, and suffix 1 and

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the "plus" sign to the extraordinary wave. Furthermore, N is the electron concentration, z_0 is the distance from the earth's surface, r_0 is the true distance from source to receiver, R_0 the earth's radius, and χ is the zenith angle of the satellite (see Fig.1). These two path lengths differ from the true distance r_0 by the same amount $\delta_{1,2}$. The above expressions can be used in a method whereby the electron concentration is determined by measuring the angle between the planes of polarization and the difference between the group path lengths on two frequencies. The combination of these two measurements is suggested as a possible approach to the measurement of the electron concentration in interplanetary space with the aid of cosmic rockets. To measure the electron concentration in interplanetary space it is necessary to have signals on frequencies $\omega_1, \omega_2, \omega_3$ which are modulated at a low frequency Ω . The close frequencies ω_1 and ω_2 can be used to measure the Faraday effect and hence the contribution to

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$\int_{r_0}^{\infty} N dr$

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due to the ionosphere, and the distant frequencies ω_1 and ω_3 to measure the difference in the group path lengths. In order that the contribution due to interplanetary space should be comparable to that due to the ionosphere, the rocket must be at a distance of 10^6 km from the earth. The reception of signals from such distances is difficult because of the low power of the transmitters on rockets. This difficulty can easily be avoided by the use of a sinusoidally modulated signal.

Acknowledgments are expressed to G.G. Getmantsev and V.L. Ginzburg for valuable advice.

There are 1 figure and 14 references: 6 Soviet and 8 non-Soviet.

ASSOCIATION: Nauchno-issledovatel'skiy radiofizicheskiy institut pri Gor'kovskom universitete

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(Scientific Research Radiophysics Institute of the Gor'kiy University)

SUBMITTED:

April 2, 1960

21166

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99/100 (also 1041, 1046)

AUTHORS: Benediktov, Ye.A., Korobkov, Yu.S. Mityakov, N.A.,
Rapoport, V.O., and Khodaleva, L.N.

TITLE: Results of Measurements of the Absorption of Radio
Waves in the Ionosphere

PERIODICAL: Izvestiya vysshikh uchebnykh zavedeniy Radiofizika,
1960, Vol.3, No.6, pp. 957-968

TEXT: Results obtained at Gor'kiy in 1959 are reported.
The total absorption in the ionosphere was measured with the aid
of the "method of two frequencies". The method is described as
follows. Suppose that the cosmic radio emission is received
simultaneously on two frequencies, f_1 and f_2 where $f_2 > f_1$.
For each of these frequencies the integral absorption of radio
waves in the ionosphere is given by:

$$\Gamma_1 = \ln(I_{01}/I_1), \quad (1)$$

where I_{01} and I_1 are the intensities of cosmic radio emission
of frequency f_1 before and after passage through the

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Ionosphere

ionosphere. If $(2\pi f_1)^2 \gg \nu^2$ and $f_1^2 \gg f_c^2$, where ν is the
effective number of collisions of electrons with ions and
neutral molecules, and f_c is the critical frequency of the
F-layer, then the integral absorption is given by:

$$\Gamma_1 = \frac{e^2}{\pi m c f_1^2} \int_0^z N \nu dz \quad (2)$$

In this expression N is the electron concentration, z is the
thickness of the absorbing layer, e and m are the charge and
mass of the electron, and c is the velocity of light. It then
follows that $\Gamma_1/\Gamma_2 = (f_2/f_1)^2$ and hence, finally, the integral
absorption for each of the frequencies is given by

$$\Gamma_1 = \frac{\ln(I_{02}/I_{01}) - \ln(I_2/I_1)}{1 - f_1^2/f_2^2} \quad (3a)$$

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and $\Gamma_2 = \Gamma_1 (f_1/f_2)^2$ (3b)

If I_{02}/I_{01} does not depend on the galactic coordinates then changes in Γ_1 with time depend only on the ratio of the two frequencies. In fact, the above intensity ratio is not independent of the galactic coordinates but this fact should not lead to large errors in the absorption measurements. Published data on the absorption of radio waves in the ionosphere during night hours shows that the absorption is frequently negligible. If the intensity ratio I_{02}/I_{01} is determined for these hours, then the absorption for any other time can be calculated from Eq. (3). It may be shown that the optimum frequency range for the above method differs from the standard method (described by Blum et al. in Ref.2 and Mitra and Shain in Ref.3) in that it does not require highly specialized apparatus or prolonged observations. The present authors have used the above method between August and

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December 1959 on 8.6 and 25 Mc/s. The results obtained show that the absorption has a characteristic maximum at noon each day, and a minimum at about 4 hrs. In August and September there is also an additional evening maximum at about 20 hrs. The magnitude of the noon maximum was found to be 1.1 db in August, 1.15 db in September, 1.2 db in October and November, and 1.6 db in December (on 18.6 Mc/s throughout). Fig.4 shows the diurnal dependence of the total absorption (continuous curve) and the absorption in the lower layers of the ionosphere (dotted curve) averaged over the periods 23rd to 31st October (Fig.4a) and 12th to 15th November (Fig.4b). The results obtained by the Radio Astronomical methods were checked by means of the pulse method described by Pigott et al. (Ref.9). Fig.5 shows the dependence of the absorption in the F-layer on the critical frequencies of the F-layer (18.5 Mc/s) (curve I - 12th to 15th November; curve II - 20th to 31st October; curve III - data from Ref.3). Acknowledgements are expressed to G.G. Getmantsev and V.L. Ginzburg for interest and advice.

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052/E114
Results of Measurements....
There are 5 figures and 13 references: 5 Soviet and 8 non-Soviet.
ASSOCIATION: Nauchno-issledovatel'skiy radiofizicheskiy institut
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Radiophysics Institute of the Gor'kiy University)
SUBMITTED: May 10, 1960

Fig. 4a

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Results of Measurements of the Absorption of Radio Waves in the Ionosphere

Fig. 46

0.5

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Results of Measurements of the Absorption of Radio Waves in the Ionosphere

Fig.5

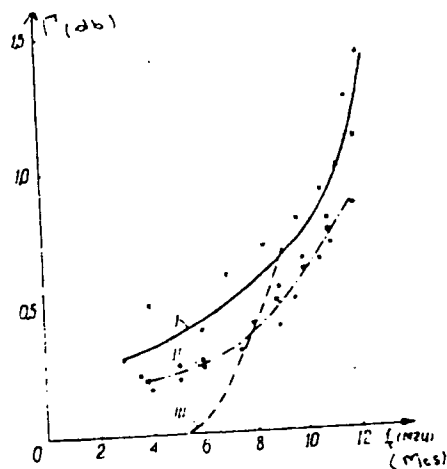


Рис. 5 Зависимость поглощения
в слое F от критических частот
слоя F (частота 18,6 мГц): I для
12-15 ноября, II для 20-31 ок-
тября, III на работе [1]

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S/141/61/004/001/003/022

EO32/E314

9.9100

AUTHORS: Benediktov, Ye.A. and Mityakov, N.A.

TITLE: On the Absorption of Cosmic Radio Emission in the Ionosphere

PERIODICAL: Izvestiya vysshikh uchebnykh zavedeniy, Radiofizika, 1961, Vol. 4, No. 1, pp. 44 - 48

TEXT: When radio waves are incident normally on the ionosphere, then for frequencies much greater than the critical frequency the absorption of these waves in the ionosphere is given by (Ginzburg - Ref. 5)

$$I' = 4,34 \frac{e^2}{\pi m c f^2} \int_0^{\infty} N_z dz = 1,16 \cdot 10^{-2} f^{-2} \int_0^{\infty} N_z dz \quad (\text{d}\theta), \quad (1)$$

where e and m is the charge and mass of the electron,
 c is the velocity of light,
 f is the frequency,

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N is the electron concentration and
ν is the effective collision frequency..

It is known (Ref. 5) that the effective collision frequency ν is determined by collisions with neutral molecules up to 150 km, while in the F-layer it is determined by collisions with ions. The magnitude of Nν can be estimated from known values of N and ν in the lower ionosphere (Nicolet - Ref. 6, Kane - Ref. 7 and Nertney - Ref. 8). These data are given in Table 1. Numerical estimates of absorption using Eq. (1) and the data in Table 1 show that the absorption in the lower layers of the ionosphere on 18.6 Mc/s is 0.3 - 0.5 db at mid-day, which is in agreement with the experimental data reported by Benediktov et al (Ref. 3). In the F-layer, the effective collision frequency is given by (Ref. 5)

$$\nu = \frac{5.5N}{T^{1/2}} \ln \left(220 \frac{T}{N^{1/2}} \right), \quad (2)$$

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On the Absorption

where T is the electron temperature. In approximate calculations it may be assumed that $T = 1\,000^\circ\text{K}$ and $N = 10^6$ and hence one obtains the approximate expression

$$\nu = \frac{45N}{T^{3/2}} \quad (2a) .$$

Substituting Eq. (2a) into Eq. (1), we have the following expression for the absorption in the F-layer

$$\Gamma_F = 0.52 f^{-2} \int \frac{N^2}{T^{3/2}} dz \quad (3) .$$

The electron concentration N is a maximum at 300 km while the temperature in the F-layer increases monotonically with

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On the Absorption

height. Since N^2 changes with height much more rapidly than $T^{-3/2}$, the above expression can be approximated by

$$\Gamma_F \approx 0.52 f^{-2} T_0^{-3/2} \int N^2 dz \quad (3a)$$

where $T_0^{-3/2}$ is an average value of $T^{-3/2}$. The electron concentration in the F-layer on the first approximation is given by (Al'pert - Ref. 9)

$$N = \begin{cases} N_0 \left(1 - \frac{z^2}{z_1^2}\right) & (z < 0) \\ N_0 \exp\left(-\frac{z}{z_1}\right) & (z > 0) \end{cases} \quad (4)$$

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where N_0 is the electron concentration in the maximum of the layer. It then follows that

$$\int_{-z_1}^{\bar{z}} N^2 dz = \frac{N_0^2}{2} \left(z_2 + \frac{16}{15} z_1 \right) \quad (5)$$

and hence

$$V_F = 0,26 f \cdot T_0 \cdot N_0^2 \left(z_2 + \frac{16}{15} z_1 \right). \quad (6)$$

Assuming standard values for the F-layer ($N_0 = 10^6$, $f_c \sim 9$ Mc/s, $z_1 = 150$ km, $z_2 = 300$ km and $T_0 = 1\,000$ °K), we find that at $f = 18.6$ Mc/s, the absorption $\Gamma_F = 1.1$ db. This is also in agreement with experimental data reported in Ref. 3. Thus, the integral absorption of radio waves in the ionosphere on frequencies considerably in excess of the critical frequency is

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largely determined by absorption in the F-layer. Eq. (6)
can also be used to determine the temperature T_o near the
maximum of the F-layer. Assuming that $N_o = 1.24 \times 10^{-8} f_c^2$,
it is found that

$$\int Ndz = 1.24 \times 10^{-8} f_c^2 z_{\text{eff}}^2 \quad (7)$$

where the effective thickness of the atmosphere is given by

$$z_{\text{eff}} = z_2 + \frac{2}{3} z_1 \quad (8)$$

Using Eqs. (6)-(8), one finds that

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$$T_0 = \left(\frac{3.1 \cdot 10^8}{\int N dz + 5 \cdot 10^{-9} z f_c^2} \frac{f^2}{f_c^2} \Gamma_F \right)^{-1/2} \quad (9)$$

Using the experimental data (Ref. 3) on the absorption in the F-layer on 18.6 Mc/s, one can calculate the product $T_0^{-3/2} z_{\phi\phi}$. In fact, using Eqs. (6)-(8), it turns out that

$$T_0^{-3/2} z_{\phi\phi} \left(1 + 0.4 \frac{z_1}{z_{\phi\phi}} \right) = 2.5 \cdot 10^{16} \frac{f^2}{f_c^2} \Gamma_F \quad (11)$$

Usually, the second factor in the brackets in Eq. (11) can be neglected. Table 2 gives the various parameters for October, 1959, as calculated from the above formulae. It is pointed out that simultaneous measurement of absorption in the F-layer and the total electron concentration $\int N dz$ can provide

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reliable information on the temperature near the maximum of the F-layer and its variation with time. Acknowledgments are expressed to V.L. Ginzburg and F.F. Getmantsev for their advice and interest. There are 2 tables and 9 references: 4 Soviet and 5 non-Soviet.

ASSOCIATION: Nauchno-issledovatel'skiy radiofizicheskiy institut pri Gor'kovskom universitete
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E032/E114

AUTHORS Vodeneyeva, D K., and Mityakov, N A
TITLE Results of experimental studies of the triple-
splitting effect in the F-layer of the ionosphere
PERIODICAL Izvestiya vysshikh uchebnykh zavedeniy
Radiofizika v 4, no. 6, 1961 1013 1019

TEXT The anisotropy of the ionosphere is usually responsible for the appearance of at least two branches on the F-layer ionograms and these are due to the ordinary and the extraordinary waves. Frequently, however, one observes an additional branch which is referred to as the z-component or the triple-splitting effect. The present authors report results of experimental studies of the latter effect in the F-layer. The observations were carried out in March 1961 at Gor'kiy. The results are in complete agreement with those reported by G R Ellis (Ref 3: J. Atm Terr Phys v 3, 263 (1953), v 8 43 (1956)). It is established that the reason for the appearance of the z-component is the interaction between

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Results of experimental studies

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obliquely incident radio waves with the radio waves which are back scattered by ionospheric irregularities. Acknowledgments are expressed to V.L. Ginzburg, G.G. Getmantsev, L.A. Skrebkova and V.O. Rapoport for their assistance in this work. N.G. Denisov is mentioned in the article.

There are 3 figures and 7 references. 3 Soviet-bloc and 4 non-Soviet-bloc. The 4 English language references read as follows Ref. 3: in text above.

Ref. 2: G.C.W. Scott: J. Geophys. Res., v. 55, 64 (1950)

Ref. 6: B. Landmark: J. Atm. Terr. Phys., v. 2, 254 (1952)

Ref. 7: R. Saranaravana: J. Atm. Terr. Phys., v. 13, 201 (1959)

ASSOCIATION Nauchno issledovatel'skiy radiofizicheskiy institut pri Gor'kovskom universitete
(Scientific Research Radiophysics Institute at Gor'kiy University)

SUBMITTED June 8, 1961

Can 1 2/2

MITYAKOV, N.A.; RAPOPORT, V.O.

Possibility for measuring the electron concentration in the upper ionosphere and in interplanetary space on the basis of plasma wave radiation. Izv. vys. ucheb. zav; radiofiz. 5 no.3:464-467 '62. (MIRA 15:7)

1. Nauchno-issledovatel'skiy radiofizicheskiy institut pri Gor'kovskom universitete.
(Electrons) (Ionosphere) (Outer space)

MITYAKOV, N. A.

Field frequency of an oscillator moving in an anisotropic
medium. Izv. vys. ucheb. zav.; radiofiz. 5 no.5:892-896 '62.
(MIRA 15:10)

1. Nauchno-issledovatel'skiy radiofizicheskiy institut pri
Gor'kovskom universitete.

(Oscillators, Electric) (Plasma(Ionized gases))

MITYAKOV, N.A.; MITYAKOV, E.Ye.; CHEREPOVITSKIY, V.A.

Results of radio observations from the artificial satellites
"Kosmos 1" and "Kosmos 2" in the Crimea. Geomag. i aer. 3 no.
5:816-822 S-O '63. (MIRA 16:11)

1. Radiofizicheskiy institut pri Gor'kovskom gosudarstvennom
universitete.

MITYAKOV, N.A.; MITYAKOVA, E.Ye.

Methodology of studying the ionospheric structure by way of
ground reception of radio signals. ~~CONFIDENTIAL~~. Geo-
mag. i aer. 3 no. 5:858-867 S-O '63. (MIRA 16:11)

1. Radiofizicheskiy institut pri Gor'kovskom gosudarstvennom
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MITYAKOV, R.A.

Letter to the editor. Izv. vpr. u sb. prav.: radiofiz. 1974.
10p. (No subject headings)

L 6846-65 EWT(a)/EWT(1)/EEC(k)-2/ENG(r)/FOC/EEC-1/EEC(t)/EEC(c)-2/EWA(h)/FSS-2
 Pb-1/Pa-1/Pa-4/Pa-5/Pi-4/Pa-4/Pg-4/Pae-2/Pt-10/Pi-4/Pl-4 AFETR/ASD(m)-3/SSD/
 AFMD/AFEM(a)/ASD(a)-5/APTGA/AFWI/AS(mf)-2/ASD(p)-3/ESD/ESD(dr)/ESD(c)/ESD(t)/BAEM(t)
 8/0141/64/007/003/0556/0559
 ACCESSION NR: AP4044111 129
 128

AUTHORS: Verukhimov, L. M.; Mityakov, N. A.

TITLE: Some methods of ionosphere research by diversity reception
of signals from artificial earth satellites

SOURCE: IVUZ. Radiofizika, v. 7, no. 3, 1964, 556-559

TOPIC TAGS: ionospheric radio wave, diversity reception, artificial
 earth satellite, electron concentration, phase measurement

ABSTRACT: The authors discuss briefly some methods for determining
 ionospheric parameters, based on diversity reception of signals
 from artificial satellites. These methods are based on calculations
 of the optical path length of the radio wave in the ionosphere, with
 allowance for Snell's law and with corrections for the horizontal
 gradients of the electron concentration. The use of the Faraday
 method in conjunction with diversity reception yields the horizontal

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components of the phase difference between the ordinary and extraordinary wave and consequently the integral gradients of the electron concentration in the direction of the line joining the two antennas. Measurements of the phase difference of coherent frequencies at two separate points makes it possible to determine directly the regular and irregular horizontal gradients of the electron concentration in a direction perpendicular to the plane of incidence of the wave. Estimates show that the necessary measurement base at 20 and 90 Mcs should be of the order of several hundred meters. It is noted in a postscript that experimental data obtained by diversity reception of signals from the satellite "Kosmos-1" disclose the presence of large inhomogeneities (larger than 500 km) with horizontal gradients exceeding 10^5 electron/cm², causing appreciable shifts in the Faraday fadings even when the antennas are separated by a distance on the order of 1 km. Orig. art. has: 10 formulas and 1 figure.

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ASSOCIATION: Nauchno issledovatel'skiy radiofizicheskiy institut
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NR REF SOV: 005

OTHER: 000

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ACCESSION NR: AT5023576 AST/TT/GS/GW

UR/0000/65/000/000/0147/0150

AUTHOR: Yerukhimov, L. M.; Mityakov, N. A.; Mityakova, E. Ye.

69

B+1

TITLE: Investigation of the ionosphere by the method of ground reception of radio signals from artificial earth satellites

SOURCE: Vsesoyuznaya konferentsiya po fizike kosmicheskogo prostranstva. Moscow, 1965. Issledovaniya kosmicheskogo prostranstva (Space research); trudy konferentsii. Moscow, Izd-vo Nauka, 1965, 147-150

TOPIC TAGS: ionosphere, ionospheric inhomogeneity, electron density, artificial satellite observation

ABSTRACT: A summary of research on the regular ionospheric structure, large-scale inhomogeneities of electron concentration, and small-scale ionospheric inhomogeneities is presented. The research in question has been conducted since 1961 using artificial earth satellites (Elektron-1 included). The regular structure of the ionosphere was studied by the measurement of the phase difference of coherent frequency signals (20—90 Mc) and the Faraday fading of 20-Mc signals from the satellites. According to data from Cosmos-1, Cosmos-2, and Explorer-7, electron concentration

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12

L 2965-66

ACCESSION NR: AT5023576

as a function of the altitude of the satellite and the time of day was in the range $0.1-1.7 \times 10^{13}$ el/cm². The corresponding average value for the exponent index k was 6.2×10^{-3} /km. The index was determined from comparisons with vertical probing data under the assumption that the electron concentration above the F layer behaves exponentially. The measurements of large-scale inhomogeneities indicated that their dimensions range from a few kilometers to a few hundred kilometers. The gradient

$$\Delta \int \frac{\partial N}{\partial s} ds \approx 10^4 \text{ el/cm}^3$$

is independent of the nonuniformity dimension l for $l > 100$ km. For $l < 100$ km, this gradient increases with l . Small-scale inhomogeneities were determined from the fluctuation of signals received at three spatially dispersed antennas. It was established that they have a clearly expressed daily course, with the maximum occurring at night. They were observed primarily at 250-350 km and ranged in size from 1 to 2 km. Orig. art. has 1 formula. [BD]

ASSOCIATION: none

SUBMITTED: 02Sep65

NO REF SOV: 008

Card 2/2 RYK

ENCL: 00

OTHER: 001

SUB CODE: ES EC

ATD PRESS: 4109

L 5315-66 EWT(d)/FBD/FSS-2/EWT(1)/FS(v)-3/EEC(k)-2/EWA(d) AST/TT/RB/GS/GW/WS-2
 UR/0000/65/000/000/0581/0606
 106
 ACCESSION NR: AT5023642

AUTHORS: Benediktov, Ye. A.; Getmantsev, G. G.; Mityakov, N. A.; Rapoport, V. O.;
Sazonov, Yu. A.; Tarasov, A. F.
 55 55 85

TITLE: Results of the intensity measurements of radio-frequency radiation at frequencies of 725 and 1525 kc by means of the apparatus installed in the satellite Elektron-2

SOURCE: Vsesoyuznaya konferentsiya po fizike kosmicheskogo prostranstva. Moscow, 1965, Issledovaniya kosmicheskogo prostranstva (Space research); trudy konferentsii. Moscow, Izd-vo Nauka, 1965, 581-606

TOPIC TAGS: artificial earth satellite, radio emission, ionosphere, atmospheric radiation, radio receiver, geomagnetic field

ABSTRACT: The results of radio-frequency measurements taken by the Elektron-2 satellite are analyzed and the equipment used is described. Two fixed-frequency receivers tuned to 725 and 1525 kc were used with a common dipole antenna. One side of the antenna was a 3.75-m metal stub, and the other side was the body of the satellite; the radiation resistance was 0.033 ohm for 725 kc and 0.146 ohm for 1525 kc for a capacitance of 46 pF. The receivers used straight amplification with 3 rf

Card 1/5

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L 5315-66

ACCESSION NR: AT5023642

21

stages and 2 af stages. The error in the absolute value of the intensity of cosmic radio emission was $\pm 30\%$ for 1525 kc and $(+30, -50)\%$ for 725 kc. The measurement results were processed by converting the output voltages to the effective temperature of radio emission. Values of effective temperature T_{eff} for a 2-hr flight near

the apogee are given in Fig. 1 on the Enclosure, where the points correspond to 1525 kc and the crosses to 725 kc. All of the data on the spectrum of cosmic radio emission indicate that for $f \leq 3-5$ Mc its intensity decreases with frequency. The profile of the electron concentration in the ionosphere was determined from its effect on radiation resistance and capacitance of the antenna. A graph of electron concentration N versus altitude h is shown in Fig. 2 on the Enclosure. Sporadic radio emission from the earth's atmosphere considerably exceeding the cosmic radio emission in intensity was recorded at both frequencies. A correlation between radio emission and the intensity of soft-electron flux is found. The distribution of radio emission indicates that electron fluxes penetrate the ionosphere primarily at latitudes of 30-50°. The authors thank Yu. V. Abramov, A. A. Andronov, B. N. Boykin, V. L. Ginsburg, V. V. Zheleznyakov, V. S. Karavanov, Yu. I. Logachev, G. A. Skuridin, and V. Yu. Trakhtengerts for aid in preparing the experiment and discussion of the results. Orig. art. has: 14 graphs, 1 diagram, 1 chart, 3 tables, and 11 formulas.

Card 2/5

L 5315-66

ACCESSION NR: AT5023642

ASSOCIATION: Vsesoyuznaya konferentsiya po fizike kosmicheskogo prostranstva,
Moscow (All-Union Conference on Space Physics)

SUBMITTED: 02Sep65

ENCL: 02

SUB CODE: ES, NP

NO REF SOV: 011

OTSER: 007

Card 3/5

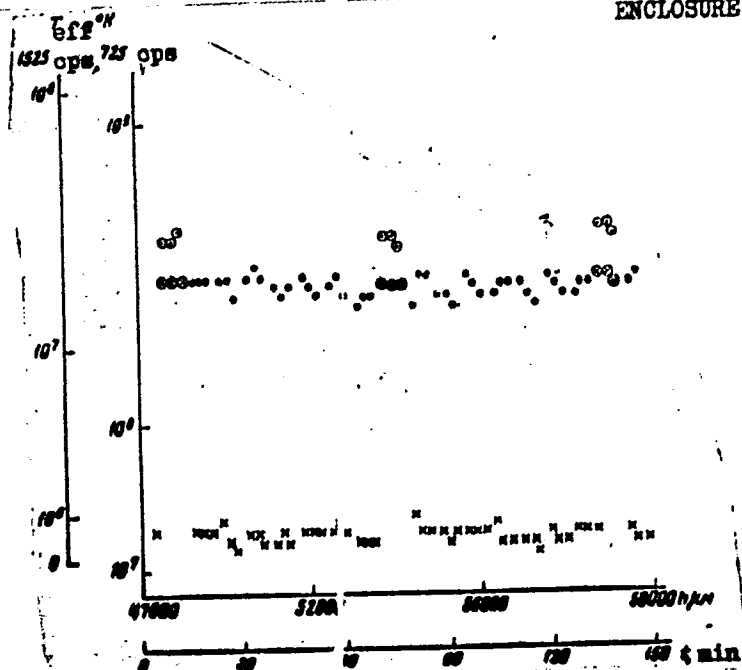
L 5315-66

ACCESSION NR: AT5023642

ENCLOSURE:

81

Fig. 1. Effective temperature versus time



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L 5315-66

ACCESSION NR: AT5023642

ENCLOSURE: 02

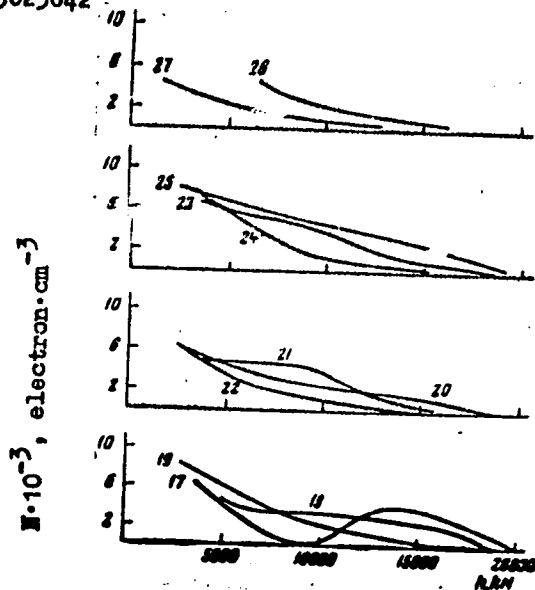


Fig. 2. Electron concentration versus altitude

GC
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L 65295-65 EWT(d)/EWT(1)/FS(v)-3/FS3-2 TT/AST/GW

ACCESSION NR: AP5021255

UR/0293/65/003/004/0618/0629
629.195.2:621.99

AUTHORS: Getmantsev, G. G.⁴⁴; Kalashnikov, M. I.⁴⁴; Bykov, V. L.⁴⁴; Benadiktov, Ya. A.⁴⁴
Yerukhinov, E. M.⁴⁴; Belikov, V. V.⁴⁴; Bakhtin, V. M.⁴⁴; Kantor, L. Ya.⁴⁴; Korobkov, S.⁴⁴
Yu, S.⁴⁴; Kunitov, M. V.⁴⁴; Mitvakov, N. A.⁴⁴; Puzirev, I. M.⁴⁴; Rapoport, V. O.⁴⁴; Sigalov,
A. G.⁴⁴; Cherepovitskiy, V. A.⁴⁴; Akin, E. A.⁴⁴

TITLE: The results of an experiment on radio communications via "Echo 2" and the moon at a frequency of 162.4 megacycles between the observatories of Jodrell Bank and Zimenki

SOURCE: Kosmicheskiye issledovaniya, v. 3, no. 4, 1965, 618-629

TOPIC TAGS: moon, satellite communication, radio telescope, radio transmission, satellite tracking, scientific research coordination / Jodrell Bank radio telescope, Zimenki observatory radio telescope, BESM 2 electronic computer

ABSTRACT: During February-March 1964 the Academy of Sciences of the USSR, NASA of the USA, and the General Post Office Department of Great Britain conducted an experiment to establish one-way radio communication at 162.4 megacycles via the passive satellite "Echo-2" and the moon. Echo-2 was used for 34 communication

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L 65295-65

ACCESSION NR: AP5021255

tests of 10-15 minutes (the time interval permitted by Echo's orbit), and the moon was used for 15 test runs between the Echo tests. The transmitting equipment at Jodrell Bank and the receiving unit of the Zimenki Observatory are described in detail. Echo orbit information furnished by NASA, visual observations, and radio tracking data from fixed stations were fed to a BESM-2 electronic calculator which provided programmed tracking control. The received signal exhibited strong fluctuations separable into two periods: 1) a 1-2 minute fluctuation associated with Echo-2 distortion from a sphere and with tracking errors; 2) a 3-10 second period associated with small surface irregularities. The rapid fluctuations varied with each test. Voice signals, slowed by a factor of 8, were barely intelligible. Telegraph, teletype, and photofacsimile transmission, in general, were unsatisfactory, but in periods of high signal-to-noise ratios intelligible messages were received. The moon transmissions were not as clear but did furnish scientific information. Unexpected transmission losses included 3-5 db for polarization losses and 1-2 db for unknown causes. The international cooperation was excellent, with the Soviet submitting a complete report. Offers for further cooperation have been extended. Orig. art. has: 3 tables, 7 figures, and 4 formulas.

ASSOCIATION: none

SUBMITTED: 18Apr65

NO REF SOV: 000

Card 2/2

ENCL: 00

OTHER: 002

SUB CODE: AA, EC

I 23430-66 EWT(d)/FSS-2/EWT(1)/EEC(k)-2/FIC/EWA(d)/EWA(h) AST/TT/GW

ACC NR: AP6012830

SOURCE CODE: UR/0293/66/004/002/0249/0256

AUTHOR: Mityakov, N. A.; Mityakova, E. Ye.; Cherepovitskiy, V. A.

ORG: none

TITLE: Results of a study of the distribution of electron concentration in the ionosphere by a method of ground reception of radio signals from Electron-1

SOURCE: Kosmicheskiye issledovaniya, v. 4, no. 2, 1966, 249-256

TOPIC TAGS: ionosphere, ionospheric electron concentration/Electron 1

ABSTRACT: The total electron concentration in the ionosphere above the maximum of the F layer was determined from ground reception of signals of Electron-1 transmitted at 20.005 and 30.0075 Mc. Observations were made during February-March 1964 at Gorky and in the Crimea with equipment capable of recording the phase difference of coherent-frequency signals. Standard PKCh-3 equipment, described earlier by Ya. L. Al'pert et al., was employed in the Crimea, while special equipment capable of recording signal amplitudes and phase differences at coherent frequencies of 20, 30, 40, and 90 Mc was developed for use at Gorky. Standard R-250 M receivers were employed. Signals from a coherent reference heterodyne were also fed to the receivers. In the presence of satellite signals, low-frequency beats were generated at the output of the receivers. After passing through narrow-band filters, the low-frequency signals were fed to a phase meter, where they were brought to a single frequency of 9 kc. On the basis of recorded phase differences, total electron concentration was

Card 1/3 UDC: 350.388.1

L 23430-66

ACC NR: AP6012830

determined to altitude z_c of the satellite from the following formula:

$$N_{no} = \int_0^{z_c} N dz$$

where N_{nc} is the vertical profile of the ionosphere passing through a point at which radio beams intersect with the maximum of the F layer. Curves showing the diurnal variation of N_{nc} for various intervals of geographic latitudes are given in the

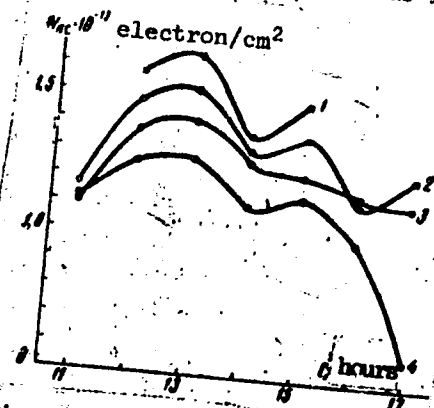


Fig. 1. Diurnal variation of the electron concentration for various geographic latitudes

1 - 51-53°; 2 - 53-55°; 3 - 55-57°; 4 - 57 to 60°.

figure. The total electron concentration was found to increase in the southward

Card 2/3

L 23430-66

ACC NR: AP6012830

direction. In conclusion, the authors avail themselves of the opportunity to thank T. I. Makarov and S. K. Malyshev for their participation in the development and preparation of the equipment; L. M. Barsukov, V. A. Vasin, and L. I. Grekov for their assistance in processing the material; and L. V. Piskunov and A. V. Potemkin for computing the ephemerides of the satellites. Orig. art. has: 6 figures and 1 table.

12

[JR]

SUB CODE: 04, 17/ SUBM DATE: 05Jun65/ ORIG REF: 006/ OTH REF: 005/ ATD PRESS:

4236

Card 3/3 *Ida*

MITYAKOV, V.I.

Device for dumping the bucket of the UEM-5 crane in conveying
concrete mix. Rats. 1 izobr. predl. v stroi. no. 75:7-8 '53.
(Concrete) (Cranes, derricks, etc.) (MIRA 7:7)

PA 43/49T48

MITYAKOV, V. S.

Oct 49

USSR/Engineering
Steam Boilers
Scale Removal

"Experiment in Washing Fire-Tube Steam Boilers
With Hydrochloric Acid," V. S. Mit'akov, Engr, "17

"Ze Ekonomiya Topliva" Vol V, No 10

Describes experiment in removing boiler scales
from steam boilers with weak solutions of hydro-
chloric acid. Method should be used only once or
twice on any one boiler.

43/49T48

1944, 1945, 1946.

1947. Vnutrikotlovoy vo sverkhvostokakh SSSR. M.: Voenizdat, 1947, No. 1, p. 1-10.

1948. Vnutrikotlovoy vo sverkhvostokakh SSSR. M.: Voenizdat, 1948, No. 1, p. 1-10.

GRINBOYM, M.Ya.; GUTOROV, V.G., ZHILYAYEV, A.V., KASATKIN, V.N.; LEVIN, P.V. [deceased], MITYAKOV, V.S., OKOROKOV, A.A.; USHAKOV, P.N.; BURKOV, G.A., laureat Stalinskoy premii, redaktor [deceased]; AYZENSHTAT, I.I., redaktor. FRIDKIN, A.M., tekhnicheskii redaktor.

[Handbook on boiler inspection] Spravochnik po kotlonadзору.
Izd. 2-e, perer. Pod obshchei red. G.A.Burkova. Moskva, Gcs.
energ. izd-vo, 1954. 568 p. [Microfilm] (MLRA 8-2)
(Boilers--Inspection)

L 8950-65 EWT(d)/FSS-2/E.T(1)/EEC(k)-2/ENG(v)/FCC/EEC-L/EEC(t)/EEC(c)-2/EWA(h)
 Pn-4, Po-4, Ps-5/Pq-4/Pac-4/Pg-4/Pae-2/Peb/Pi-4/Pk-4/Pl-4/Pb-4 AFMD(t)/RAEM(t)/
 ACCESSION NR: AP4043250 ASD(m)-3/AS(mp)-2/ S/0203/64/004/004/0668/0674
 AFETR/AFMDC/ESD(t)/ASD(p)-3/ESD(c)/SSD/ASD(a)-5/AFTC(a)
 AUTHOR: Mityakova, E. Ye. AST/GW/WS

TITLE: ^{12/}Measurements of electron concentration in the ionosphere based on observations of the Faraday effect of the radio signals of artificial earth satellites

SOURCE: Geomagnetizm i aeronomiya, v. 4, no. 4, 1964, 668-674

TOPIC TAGS: ionospheric electron concentration, electron concentration, Faraday effect, artificial earth satellite, artificial earth satellite signal, satellite signal

ABSTRACT: It is suggested that one of the possible methods of determining electron concentrations in the ionosphere is the method based on amplitude recordings of radio signals received from artificial earth satellites. In this connection, methods of processing the Faraday fadings of these signals in order to determine electron concentrations for a vertical column of an even cross section are established, and the respective formulas and the limits

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ACCESSION NR: AP4043250

of their application are derived. The method is applied to signals received from Cosmos-1 and Cosmos-2 in the Crimea (March-April 1962) and from Explorer-VII in Gorky (beginning of 1962). It is concluded that although the processing involving Faraday-effect recordings is not as accurate as the processing which involves phase recording, the amplitude method has, nevertheless, the advantage of simplifying the experiment which takes place in the course of regular observations. Orig. art. has: 4 figures, 3 tables, and 17 formulas.

ASSOCIATION: Radiofizicheskiy institut, Gor'kovskiy gosudarstvennyy universitet (Institute of Radio Physics, Gorky State University)

SUBMITTED: 11Dec63

ATD PRESS: 3105

ENCL: 00

SUB CODE: AA, EM

NO REF SOV: 003

OTHR: 008

Card 2/2

9.9120 (also 1041, 1046)
9.9400

S/141/60/003/006/004/025
E032/E114

AUTHORS: Mityakova, E.Ye., Mityakov, N.A., and Rapoport, V.O.

TITLE: On the Measurement of the Electron Concentration in the Ionosphere and in Interplanetary Space

PERIODICAL: Izvestiya vysshikh uchebnykh zavedeniy, Radiofizika, 1960, Vol.3, No.6, pp. 949-956

TEXT: A brief review is given of the available methods for the determination of the electron concentration in the ionosphere with the aid of artificial earth satellites. Using the quasi-longitudinal approximation, an expression is obtained for the phase and group paths for a signal emitted from an artificial earth satellite towards a spherical earth. It is shown using the results of Al'pert et al (Ref.11) that the phase path length is given by

$$n_{1,2}^2 = 1 - \frac{4\pi e^2 N}{m \omega(\omega \pm \omega_L)} = 1 - \frac{2aN}{\omega(\omega \pm \omega_L)} \quad (1)$$

and

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E032/E114

On the Measurement of the Electron Concentration in the
Ionosphere and in Interplanetary Space

$$L_{\phi 1,2} = r_0 + \frac{a}{\omega \cos \gamma} \left[\int_0^z \frac{N}{\omega \pm \omega_L} dz - \operatorname{tg}^2 \gamma \int_0^z \frac{Nz}{R_0 (\omega \pm \omega_L)} dz \right]. \quad (5)$$

and the group path length is given by

$$L_{rpl,2} = \int_A^B \frac{\partial (n_{1,2} \omega)}{\partial \omega} d\omega_{1,2} \quad (6)$$

and

$$L_{rpl,1} = r_0 + \frac{a}{\omega \cos \gamma} \left[\int_0^z \frac{N}{\omega \pm \omega_L} dz - \operatorname{tg}^2 \gamma \int_0^z \frac{Nz}{R_0 (\omega \pm \omega_L)} dz \right]. \quad (7)$$

In these expressions $\omega_L = (eH_0/mc) \cos \gamma$, γ is the angle
between the earth's magnetic field and the wave normal, suffix 2
and the "minus" sign refer to the ordinary wave, and suffix 1 and
Card 2/5

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S/141/60/003/006/004/025
E032/E114

On the Measurement of the Electron Concentration in the
Ionosphere and in Interplanetary Space

the "plus" sign to the extraordinary wave. Furthermore, N is the
electron concentration, z_0 is the distance from the earth's
surface, r_0 is the true distance from source to receiver,
 R_0 the earth's radius, and χ is the zenith angle of the satellite
(see Fig.1). These two path lengths differ from the true distance
 r_0 by the same amount $c_{1,2}$. The above expressions can be used
in a method whereby the electron concentration is determined by
measuring the angle between the planes of polarization and the
difference between the group path lengths on two frequencies. The
combination of these two measurements is suggested as a possible
approach to the measurement of the electron concentration in
interplanetary space with the aid of cosmic rockets. To measure
the electron concentration in interplanetary space it is necessary
to have signals on frequencies $\omega_1, \omega_2, \omega_3$ which are modulated at
a low frequency Ω . The close frequencies ω_1 and ω_2 can be
used to measure the Faraday effect and hence the contribution to

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$$\int_0^{r_0} N dr$$

21165

S/141/60/003/006/004/025
EO32/E114

On the Measurement of the Electron Concentration in the
Ionosphere and in Interplanetary Space

due to the ionosphere, and the distant frequencies ω_1 and ω_3 to
measure the difference in the group path lengths. In order that the
contribution due to interplanetary space should be comparable to
that due to the ionosphere, the rocket must be at a distance of
 10^6 km from the earth. The reception of signals from such
distances is difficult because of the low power of the transmitters
on rockets. This difficulty can easily be avoided by the use of
a sinusoidally modulated signal.

Acknowledgments are expressed to G.G. Getmantsev and
V.L. Ginzburg for valuable advice.

There are 1 figure and 14 references: 6 Soviet and 8 non-Soviet.

ASSOCIATION: Nauchno-issledovatel'skiy radiofizicheskiy institut
pri Gor'kovskom universitete
Card 4/5 (Scientific Research Radiophysics Institute of the
Gor'kiy University)
SUBMITTED: April 2, 1960

MITYAKOV, N.A.; MITYAKOVA, E.Ye.

Methodology of studying the ionospheric structure by way of
ground reception of radio signals ~~from the ionosphere~~. Geo-
mag. i aer. 3 no. 5:858-867 S-O '63. (MIRA 16:11)

1. Radiofizicheskiy institut pri Gor'kovskom gosudarstvennom
universitete.

L 2965-66/ EWT(d)/FSS-2/EWT(1)/FS(v)-3/EPA(sp)-2/EEC(k)-2/FSS/EWA(d)/EWA(h)

ACCESSION NR: AT5023576 AST/TT/GS/GW

UR/0000/65/000/000/0147/0150

AUTHOR: Yerukhimov, L. M.; Mityakov, N. A.; Mityakova, E. Ye.

69

B+1

TITLE: Investigation of the ionosphere by the method of ground reception of radio signals from artificial earth satellites

SOURCE: Vsesoyuznaya konferentsiya po fizike kosmicheskogo prostranstva. Moscow, 1965. Issledovaniya kosmicheskogo prostranstva (Space research); trudy konferentsii. Moscow, Izd-vo Nauka, 1965, 147-150

TOPIC TAGS: ionosphere, ionospheric inhomogeneity, electron density, artificial satellite observation

ABSTRACT: A summary of research on the regular ionospheric structure, large-scale inhomogeneities of electron concentration, and small-scale ionospheric inhomogeneities is presented. The research in question has been conducted since 1961 using artificial earth satellites (Elektron-1 included). The regular structure of the ionosphere was studied by the measurement of the phase difference of coherent frequency signals (20—90 Mc) and the Faraday fading of 20-Mc signals from the satellites. According to data from Cosmos-1, Cosmos-2, and Explorer-7, electron concentration

Card 1/2

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12

12

L 2965-66

ACCESSION NR: AT5023576

as a function of the altitude of the satellite and the time of day was in the range 0.1--1.7 x 10¹³ el/cm². The corresponding average value for the exponent index κ was 6.2 x 10⁻³/km. The index was determined from comparisons with vertical probing data under the assumption that the electron concentration above the F layer behaves exponentially. The measurements of large-scale inhomogeneities indicated that their dimensions range from a few kilometers to a few hundred kilometers. The gradient

$$\Delta \int_{x_0}^x \frac{\partial N}{\partial x} dx \approx 10^4 \text{ el/cm}^3$$

is independent of the nonuniformity dimension l for $l > 100$ km. For $l < 100$ km, this gradient increases with l . Small-scale inhomogeneities were determined from the fluctuation of signals received at three spatially dispersed antennas. It was established that they have a clearly expressed daily course, with the maximum occurring at night. They were observed primarily at 250--350 km and ranged in size from 1 to 2 km. Orig. art. has 1 formula.

[BD]

ASSOCIATION: none

SUBMITTED: 02Sep65

NO REF SOV: 008

Card 2/2 Bvk

ENCL: 00

OTHER: 001

SUB CODE: ES EC

ATD PRESS: 4/09

L 23430-66 EWT(d)/FSS-2/EWT(1)/ERC(k)-2/FCC/EWA(d)/EWA(h) ASI/II/GW

ACC NR: AP6012830

SOURCE CODE: UR/0293/66/004/002/0249/0256

AUTHOR: Mityakov, N. A.; Mityakova, E. Ye.; Cherepovitskiy, V. A.

ORG: none

TITLE: Results of a study of the distribution of electron concentration in the ionosphere by a method of ground reception of radio signals from Electron-1

SOURCE: Kosmicheskiye issledovaniya, v. 4, no. 2, 1966, 249-256

TOPIC TAGS: ionosphere, ionospheric electron concentration/Electron 1

ABSTRACT: The total electron concentration in the ionosphere above the maximum of the F layer was determined from ground reception of signals of Electron-1 transmitted at 20.005 and 30.0075 Mc. Observations were made during February-March 1964 at Gorky and in the Crimea with equipment capable of recording the phase difference of coherent-frequency signals. Standard PKCh-3 equipment, described earlier by Ya. L. Al'pert et al., was employed in the Crimea, while special equipment capable of recording signal amplitudes and phase differences at coherent frequencies of 20, 30, 40, and 90 Mc was developed for use at Gorky. Standard R-250 M receivers were employed. Signals from a coherent reference heterodyne were also fed to the receivers. In the presence of satellite signals, low-frequency beats were generated at the output of the receivers. After passing through narrow-band filters, the low-frequency signals were fed to a phase meter, where they were brought to a single frequency of 9 kc. On the basis of recorded phase differences, total electron concentration was

Card 1/3 UDC: 350.388.1

L 23430-66

ACC NR: AP6012830

determined to altitude z_c of the satellite from the following formula:

$$N_{nc} = \int_0^{z_c} N dz$$

where N_{nc} is the vertical profile of the ionosphere passing through a point at which radio beams intersect with the maximum of the F layer. Curves showing the diurnal variation of N_{nc} for various intervals of geographic latitudes are given in the

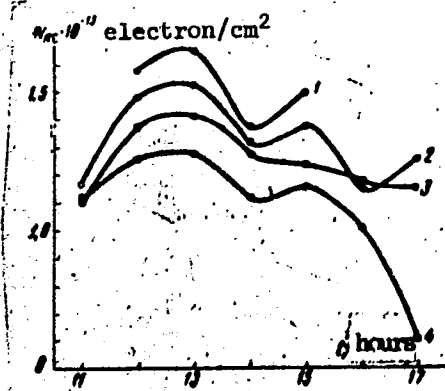


Fig. 1. Diurnal variation of the electron concentration for various geographic latitudes

1 - 51-53°; 2 - 53-55°; 3 - 55-57°; 4 - 57 to 60°.

figure. The total electron concentration was found to increase in the southward
Card 2/3

L 23430-66

ACC NR: AP6012830

direction. In conclusion, the authors avail themselves of the opportunity to thank T. I. Makarov and S. K. Malyshev for their participation in the development and preparation of the equipment; L. M. Barsukov, V. A. Vasin, and L. I. Grekov for their assistance in processing the material; and L. V. Piskunov and A. V. Potemkin for computing the ephemerides of the satellites. Orig. art. has: 6 figures and 1 table. [JR]

SUB CODE: 04, 17/ SUBM DATE: 05Jun65/ ORIG REF: 006/ OTH REF: 005/ ATD PRESS: 4236

Card

3/3 *Ida*

MAKKAVEYEV, N.I., prof.; LAPTEV, M.I.; MITYAKOVA, M.N.; KONDRAKHOVA, Ye.I.;
SHANKIN, P.A.; RZHANITSYN, N.A.; RABKOVA, Ye.K.; VYKHLOV, K.P.;
CHALOV, R.S.

[Planning the navigable channels of unregulated rivers.]
Proektirovanie sudovykh khodov na svobodnykh rekakh. Moskva,
Transport, 1964. 261 p. (Moscow. Tsentral'nyi
nauchno-issledovatel'skii institut ekonomiki i ekspluatatsii
vodnogo transporta. Trudy, no. 36). (MIRA 18:12)

BERDINSKIY, I.S.; MITZANIN, V.P.

Synthesis of unsymmetrical diphenylhydrazides of carboxylic acids.
Trudy Perm. farm. inst. no.1:159-161 '59. (MIRA 15:1)

1. Permskiy farmatsevticheskiy institut, kafedra organicheskoy i
biologicheskoy khimii. (HYDRAZIDES)

1944-1945, 1946, 1947, 1948, 1949, 1950, 1951, 1952, 1953, 1954, 1955, 1956, 1957, 1958, 1959, 1960, 1961, 1962, 1963, 1964, 1965, 1966, 1967, 1968, 1969, 1970, 1971, 1972, 1973, 1974, 1975, 1976, 1977, 1978, 1979, 1980, 1981, 1982, 1983, 1984, 1985, 1986, 1987, 1988, 1989, 1990, 1991, 1992, 1993, 1994, 1995, 1996, 1997, 1998, 1999, 2000, 2001, 2002, 2003, 2004, 2005, 2006, 2007, 2008, 2009, 2010, 2011, 2012, 2013, 2014, 2015, 2016, 2017, 2018, 2019, 2020, 2021, 2022, 2023, 2024, 2025, 2026, 2027, 2028, 2029, 2030, 2031, 2032, 2033, 2034, 2035, 2036, 2037, 2038, 2039, 2040, 2041, 2042, 2043, 2044, 2045, 2046, 2047, 2048, 2049, 2050, 2051, 2052, 2053, 2054, 2055, 2056, 2057, 2058, 2059, 2060, 2061, 2062, 2063, 2064, 2065, 2066, 2067, 2068, 2069, 2070, 2071, 2072, 2073, 2074, 2075, 2076, 2077, 2078, 2079, 2080, 2081, 2082, 2083, 2084, 2085, 2086, 2087, 2088, 2089, 2090, 2091, 2092, 2093, 2094, 2095, 2096, 2097, 2098, 2099, 2100, 2101, 2102, 2103, 2104, 2105, 2106, 2107, 2108, 2109, 2110, 2111, 2112, 2113, 2114, 2115, 2116, 2117, 2118, 2119, 2120, 2121, 2122, 2123, 2124, 2125, 2126, 2127, 2128, 2129, 2130, 2131, 2132, 2133, 2134, 2135, 2136, 2137, 2138, 2139, 2140, 2141, 2142, 2143, 2144, 2145, 2146, 2147, 2148, 2149, 2150, 2151, 2152, 2153, 2154, 2155, 2156, 2157, 2158, 2159, 2160, 2161, 2162, 2163, 2164, 2165, 2166, 2167, 2168, 2169, 2170, 2171, 2172, 2173, 2174, 2175, 2176, 2177, 2178, 2179, 2180, 2181, 2182, 2183, 2184, 2185, 2186, 2187, 2188, 2189, 2190, 2191, 2192, 2193, 2194, 2195, 2196, 2197, 2198, 2199, 2200, 2201, 2202, 2203, 2204, 2205, 2206, 2207, 2208, 2209, 2210, 2211, 2212, 2213, 2214, 2215, 2216, 2217, 2218, 2219, 2220, 2221, 2222, 2223, 2224, 2225, 2226, 2227, 2228, 2229, 2230, 2231, 2232, 2233, 2234, 2235, 2236, 2237, 2238, 2239, 2240, 2241, 2242, 2243, 2244, 2245, 2246, 2247, 2248, 2249, 2250, 2251, 2252, 2253, 2254, 2255, 2256, 2257, 2258, 2259, 2260, 2261, 2262, 2263, 2264, 2265, 2266, 2267, 2268, 2269, 2270, 2271, 2272, 2273, 2274, 2275, 2276, 2277, 2278, 2279, 2280, 2281, 2282, 2283, 2284, 2285, 2286, 2287, 2288, 2289, 2290, 2291, 2292, 2293, 2294, 2295, 2296, 2297, 2298, 2299, 2300, 2301, 2302, 2303, 2304, 2305, 2306, 2307, 2308, 2309, 2310, 2311, 2312, 2313, 2314, 2315, 2316, 2317, 2318, 2319, 2320, 2321, 2322, 2323, 2324, 2325, 2326, 2327, 2328, 2329, 2330, 2331, 2332, 2333, 2334, 2335, 2336, 2337, 2338, 2339, 2340, 2341, 2342, 2343, 2344, 2345, 2346, 2347, 2348, 2349, 2350, 2351, 2352, 2353, 2354, 2355, 2356, 2357, 2358, 2359, 2360, 2361, 2362, 2363, 2364, 2365, 2366, 2367, 2368, 2369, 2370, 2371, 2372, 2373, 2374, 2375, 2376, 2377, 2378, 2379, 2380, 2381, 2382, 2383, 2384, 2385, 2386, 2387, 2388, 2389, 2390, 2391, 2392, 2393, 2394, 2395, 2396, 2397, 2398, 2399, 2400, 2401, 2402, 2403, 2404, 2405, 2406, 2407, 2408, 2409, 2410, 2411, 2412, 2413, 2414, 2415, 2416, 2417, 2418, 2419, 2420, 2421, 2422, 2423, 2424, 2425, 2426, 2427, 2428, 2429, 2430, 2431, 2432, 2433, 2434, 2435, 2436, 2437, 2438, 2439, 2440, 2441, 2442, 2443, 2444, 2445, 2446, 2447, 2448, 2449, 2450, 2451, 2452, 2453, 2454, 2455, 2456, 2457, 2458, 2459, 2460, 2461, 2462, 2463, 2464, 2465, 2466, 2467, 2468, 2469, 2470, 2471, 2472, 2473, 2474, 2475, 2476, 2477, 2478, 2479, 2480, 2481, 2482, 2483, 2484, 2485, 2486, 2487, 2488, 2489, 2490, 2491, 2492, 2493, 2494, 2495, 2496, 2497, 2498, 2499, 2500, 2501, 2502, 2503, 2504, 2505, 2506, 2507, 2508, 2509, 2510, 2511, 2512, 2513, 2514, 2515, 2516, 2517, 2518, 2519, 2520, 2521, 2522, 2523, 2524, 2525, 2526, 2527, 2528, 2529, 2530, 2531, 2532, 2533, 2534, 2535, 2536, 2537, 2538, 2539, 2540, 2541, 2542, 2543, 2544, 2545, 2546, 2547, 2548, 2549, 2550, 2551, 2552, 2553, 2554, 2555, 2556, 2557, 2558, 2559, 2560, 2561, 2562, 2563, 2564, 2565, 2566, 2567, 2568, 2569, 2570, 2571, 2572, 2573, 2574, 2575, 2576, 2577, 2578, 2579, 2580, 2581, 2582, 2583, 2584, 2585, 2586, 2587, 2588, 2589, 2590, 2591, 2592, 2593, 2594, 2595, 2596, 2597, 2598, 2599, 2600, 2601, 2602, 2603, 2604, 2605, 2606, 2607, 2608, 2609, 2610, 2611, 2612, 2613, 2614, 2615, 2616, 2617, 2618, 2619, 2620, 2621, 2622, 2623, 2624, 2625, 262

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DATE: 10/26/2011

1. Legality of the contract. The contract is legal and enforceable under the laws of the State of New York.

MITYANINA, I.V.

The stratigraphic significance of foraminifera of Jurassic deposits
in southeastern White Russia. Paleont. i stratigr. SSR no.1:108-
173 '55. (MIRA 10:1)

(White Russia--Foraminifera, Fossil)

AUTHOR: Mityanina, I.V.

5-3-1957

TITLE: Stratigraphy of Jurassic Sediments in Belorussia According to the Study of Foraminifera Data. Stratigrafiya yurskikh otlozheniy Belorussii po dannym izucheniya foraminifer.

PERIODICAL: Byulleten' Moskovskogo Obshchestva Ispytateley Prirody, Otdel Geologicheskiiy, 1957, # 4, p. 177. USSR

ABSTRACT: The Middle- and Upper-Jurassic sediments in Belorussia are widespread in the Pripyat depression, in the western part of the Moscow depression and in the Brest depression of the L'vov syncline. The Upper-Jurassic system is represented by Callovian and Oxford sediments. The author describes various sediments and species of foraminifera characteristic of them.

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MITYANINA, I.V.

New data on Jurassic sediments of Grodno Province. Dokl. AN BSSR
(MIRA 12:10)
3 no.5:217-219 My '59.

1. Predstavleno akademikom AN BSSR I.S. Lupinovichem.
(Grodno Province--Geology, Stratigraphic)

MITYANINA, I.V. [Mitsianina, I.V.]

Stratigraphy of Jurassic sediments of White Russia.
Vestsi AN BSSR.Ser.fiz.-tekhn. no.4:104-107 '59.
(MIRA 13:4)

(White Russia--Geology, Stratigraphic)

MITYANINA, I.V.

Stratigraphic division of the Jurassic of White Russia.
Trudy VNIIGI no.29:91-95 v.2. 2, '61. (MIRA 14:7)
(White Russia—Geology, Stratigraphic)

MITYANINA, I.V.

Foraminifers of the Upper Oxfordian in White Russia. Paleont.1
stratigr. BSSR no.4:122-189 '63. (MIRA 17:4)

Mityanskiy, G.F.

AUTHOR: Mityanskiy, G.F.

109-12-3/15

TITLE: Migration of Barium on the Surface of Certain Metals
(Migratsiya bariya po poverkhnosti nekotorykh metallov)

PERIODICAL: Radiotekhnika i Elektronika, 1957, Vol.II, No.12,
pp. 1491 - 1496 (USSR).

ABSTRACT: The problem was already considered by several authors, in particular, by Drechsler (Ref.6) and Schaefer and White (Ref.7). A more extensive investigation of the problem was carried out by the author and the results are reported in this paper. The measurements were carried out by means of a cylindrical thermo-electronic projector (Ref.10). A mixture of barium beryllate and tantalum powder was used as the source of barium in the investigated tube, in which the pressure was initially reduced to 5×10^{-8} mmHg. The tube was then sealed off and the pressure reduced to about 2×10^{-9} mmHg. The cathode was in the form of a filament and it was coated with a multi-atomic layer of barium. Theoretically, the process of the migration of barium can be represented by:

$$x^2 = 4D_T t$$

Card 1/3 where x is the displacement of the boundary of the barium

109-12-3/16

Migration of Barium on the Surface of Certain Metals.

film, t is the time necessary to produce the displacement and D_T is the value of the migration coefficient at a given temperature T of the filament. If D_T is measured at several temperatures, it is possible to obtain the value of the activation energy of the migration, Q_M , and a constant D_∞ which are related by equation:

$$D_T = D_\infty e^{-\frac{Q_M}{KT}}.$$

The measurements were carried out on filaments having a length of 180 mm and a diameter of 0.1 mm and the following curves for the barium layers on pure tungsten were taken: $1/4x^2$ as a function of t (Fig.1), distribution of D_T along the filament and D_T against $1/T$ (Figs. 2 and 3). Similar measurements were made for barium layers on carburised tungsten, rhenium-coated tungsten and platinum-coated tungsten. It is concluded that the values of D_T and Q_M for the barium migration on pure tungsten,

Card2/3 carburised tungsten and rhenium-coated tungsten differ very little

109-12-3/15

Migration of Barium on the Surface of Certain Metals

over the investigated range of temperatures (940 to 1400 °K). On platinum-coated tungsten, the migration coefficient is larger and the activation energy is lower than that of the other three metals; in platinum, the migration of barium is faster than in other metals, especially at comparatively low temperatures. The author thanks Prof. N.D. Morgulis for directing this work. There are 4 figures, 1 table and 15 references, 2 of which are Slavic.

ASSOCIATION: Physics Institute of the Ac.Sc. Ukrainian SSR, Kiyev.
(Institut fiziki AN USSR, g. Kiyev)

SUBMITTED: May 8, 1957

AVAILABLE: Library of Congress

Card 3/3

ELECTRON PHYSICS

Migration of Barium on the Surface of Certain Metals' by G. F. Matyanskiy
Institute of Physics Academy of Sciences, Ukrainian SSR Kiev. Radio-
tekhnika i Elektronika, No 12 December 1993 pp 1491-1496.

Using a cylindrical thermionic projector at high vacuum (pressure less than 3×10^{-9} mm mercury), a comparative study is made of the migration of barium on the surface of pure tungsten, carbided tungsten, tungsten coated (electrolytically) by a layer of rhenium, and tungsten covered with a coating of platinum.

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S/185/60/005/001/006/018

A151/A029

9.4160 (320, 1003, 1105)
26.1512

AUTHORS: Borzyak, P G ; Marchuk, P.M.; Mityans'kiy, G.F.

TITLE: Photo-Electronic Emission of the Intermetallic Compounds Mg_2Sn and $InAs$

PERIODICAL: Ukrayins'kyy Fizychnyy Zhurnal, 1960, Vol. 5, No. 1, pp. 65 - 74

TEXT: Only the spectral characteristics of the photo-effect of the $A^{II}B^{IV}$ and $A^{III}B^V$ - type compounds are studied which are characterized by small, one-or-der widths of the restricted energy zone. The films Mg_2Sn were obtained by means of the condensation of a tungsten strip cleaned in a vacuum. For studying the film of a changeable composition of $Mg - Sn$, the distribution of the thermo-electronic effect of the yield φ_{temp} [ABSTRACTOR'S NOTE: Subscript temp (temperature) stands for the original T (temperatura)] along the strip 5 was determined. The results are given in Figure 3, where the curve 1 shows the distribution of φ_{temp} along the surface of the tungsten strip. Further, the film was applied on the strip, which again was cleaned in a vacuum. After that, the curve 2 was obtained. A repetition of this cycle yielded a curve which coincided with the curve 2. After a third investigation of the film, a distribution was obtained for it which

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Photo-Electronic Emission of the Intermetallic Compounds Mg_2Sn and $InAs$

is illustrated by the curve 3. For various sections of the film characterized by the curve 3, the authors have determined the spectral photo-electric sensitivity I_1 toward I_2 of the Cs_3Sb -photocathode: $\frac{I_1}{I_2} = (\lambda)$. The results are shown in Figure 4. Each curve is marked by a figure being the coordinate of the investigated section according to the data of Figure 3. The photo-electronic properties of Mg are characterized by the curve 28, those of Mg_2Sn by the curves 10, 13 and 16. It was established that the optimum value of the yield effect in respect to the photo-electronic emission is achieved for a metal surface at values t within the limits 4-20 and for the surface of Mg_2Sn at $t \approx 70$. It can also be seen that in the case of a photoelectric yield effect of only about 2 ev, the values of the quantum yield in Mg_2Sn remain small, reaching only the tenth part of a percent even at the highest values of $h\nu = 5$ ev. For studying the photo-electronic emission of $InAs$, an investigation was carried out of the surface of the break of a massive crystalline sample obtained in a high vacuum. The results of photo-electric measurements conducted on a clean, newly-obtained surface are shown by dots on the curve 1 in Figure 7. In the case of a consecutive deposition of BaO molecules on the surface of $InAs$, the yield effect is increasing, according to which the characteristics 2,3 are obtained. The curve 4 corresponds to a surface with a small amount of BaO .

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Photo-Electronic Emission of the Intermetallic Compounds Mg_2Sn and $InAs$

surface with a concentration of barium oxide molecules which is already higher than the optimum concentration. For comparison Figure 7 depicts also the curve obtained for a Cs_3Sb -photocathode, which at $\lambda = 400 \text{ m}\mu$ had a quantum yield of 18%. A comparison of the spectral characteristics of the 3 intermetallic compounds Mg_2Sn , $InAs$ and Cs_3Sb shows that the first two compounds differ from the latter one by efficiency values and appearance. Even at a distance of 3 eV from the border, they do not show any tendency to saturation and have lesser efficiencies than Cs_3Sb by one order or more. By making a comparison of the values $\Delta\phi = \phi - \phi_0$ for Mg_2Sn , $InAs$, and Cs_3Sb (Ref. 7) including here the data for Ge (Ref. 8) the tendency toward a decrease of $\Delta\phi$ is noted which occurs when the energy of the electronic affinity is also decreasing. There is, however, no direct proportional relationship between the electronic affinity and $\Delta\phi$, which shows that there are still other factors affecting the value $\Delta\phi$. There are 8 figures, 2 tables and references: 6 Soviet, 1 English and 1 German.

ASSOCIATION: Instytut Fizyki AN URSR (Institute of Physics, AS UkrSSR).

APPROVED FOR RELEASE: 06/14/2000

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SUBMITTED: July 6, 1959

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Photo-Electronic Emission of the Intermetallic Compounds Mg_2Sn and $InAs$

Figure 3:

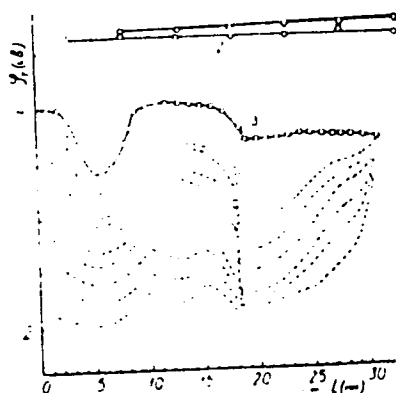


Fig. 3

Figure 4:

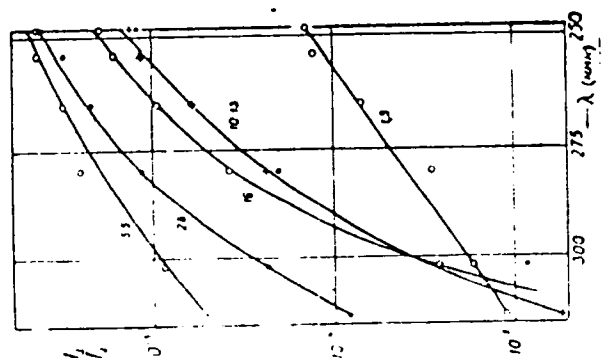


Fig. 4

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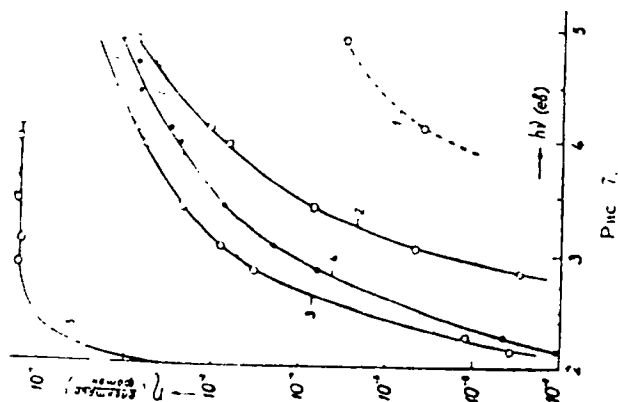
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A151/A029

Photo-Electronic Emission of the Intermetallic Compounds Mg_2Sn and $InAs$

Figure 7:



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MITYASHEV, B.N.

Noiseproof feature of one method for time discrimination of pulse signals. Nauch.dokl.vys.shkoly; radiotekh. i elektron. no.2:7-12 '58.
(MIRA 12:1)

1. Kafedra radiotekhniki Moskovskogo fiziko-tekhnicheskogo instituta.
(Radio frequency modulation---Receivers and reception)

SOV/162-58-3-4/26

9(4)

AUTHORS:

Mityashev, B.N., and Tsirlin, A.I.

TITLE:

Reducing the Influence of Sinusoidal Noise on the Pulse Signal Reception (Obshchen'shenii vliyaniya sinusoidal'noy pomekhi na priyem impul'snykh signalov)

PERIODICAL:

Nauchnyye doklady vysshey shkoly, Radiotekhnika i elektronika, 1958, Nr 3, pp 25-32 (USSR)

ABSTRACT:

The authors investigate the sine noise suppression with pulse signal reception and suggests a synchronous oscillator which produces oscillations close to the noise frequency. However, according to Ye.I. Manayev, this oscillator frequency will not be synchronous to the noise frequency. Figure 4 shows a block diagram of such a synchronous noise suppressor. A self-oscillator producing sine oscillations may be used as a synchronous generator. The block diagram of such a generator is shown by figure 6. The experimental device built according to this block diagram was somewhat bulky, containing seven vacuum tubes and

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REV/100-3-3-4/

AUTHOR: Mitchell, S. H.

AUTHOR: ME/ANOV, B. A.
TITLE: Transfer of the fluctuation noise and phase signal through a Time Discriminator (perenosdeniye fluktuatsionnoy shumy i faze signala cherez vremennyy diskriminator)

PERIODICAL: Radiotekhnika i elektronika, 1958, Vol 3, No 9, pp 1144-1157 (USSR)

[illegible]

007/10 -3-9-4/20

Transfer of the Fluctuation Noise and Puls. Signal Through a Time Discriminator

integrator. From Eqs. (37) and (38) it is seen that the mutual correlation coefficient of the system is expressed by Eq. (39). The values of this coefficient for the three types of filter are plotted in Fig. 7. If the discriminator receives pulse signals and the fluctuation noise, the analysis of the system can be carried out by considering the conditions in three distinct regions; these are indicated in Fig. 8. For this case, the dispersion of the output signal for region 1 is given by Eq. (40); in region 2 it is given by Eq. (41) or Eq. (42) while in region 3 it is given by Eq. (43) or Eq. (44). The resultant dispersion of the noise at the output is given by Eq. (45). The value of the output of the discriminator can be calculated from Eq. (46), where U_0 is the output voltage of the discriminator. The output voltage is given by

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observations. The ... and ...

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1000 /

MITYASHOV, B. N.

AUTHOR: Mityashov, B. N.,
Society member of the

TITLE: Method for the Improvement of the Demodulation Accuracy
of Pulses (Metod povysheniya tochnosti demodulyatsii
impul'sov)

PERIODICAL: Radiotekhnika, 1958, Vol. 13, Nr 5, pp. 55-62 (USSR)

ABSTRACT: Here a new method for the demodulation of pulses is shown, which was worked out by the author, and by which compared with the otherwise usual method a considerably better demodulation accuracy is reached. In the new method it is suggested to form a response (in its shape an isosceles triangle) for each signal pulse. The amplitude of these triangular pulses must be proportional to the amplitude of the signal pulses and the duration at the base must be equal to the two repetition periods. As a result of the superposition of such responses a signal is obtained in which the voltage from one amplitude value to the other changes according to a straight line. On this occasion the output signal is retarded for one repetition period with regard to the modulating function. By this the accuracy of the

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Method for the Improvement of the Demodulation Accuracy
of Pulses

demodulation is better. The reproduction accuracy is particularly high in those parts of the modulated function where the curvature is very flat. By means of low frequency filters in this method the also demodulation accuracy can be increased. The magnitude of the gain in demodulation accuracy in this method is evaluated. From the obtained tables can be seen that in case of additional connection of the filters the accuracy of the demodulation improves. The value of the relative gain in accuracy, however, remains the same as in the case when filters are not used. The data given in the tables also show the error magnitudes in the reproduction of a modulating function with pulse demodulation. Two demodulator schemes are given:

- 1) With formation of triangular pulses and
- 2) with formation of a growth function. The first type remains more complicated in spite of all simplifications than the scheme of the usually applied step-demodulator. Besides, the characteristics of both demodulator channels must be equalized. The second type is free from both deficiencies. In this the linear approximation of the

Card 2/3

В. С. Милославский
О проектной способности аппаратуры каналов связи

Ю. Н. Мартынов

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Н. А. Тихонов

Задачи теории коррелированных сигналов

Б. Н. Митин

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А. Н. Файнман

Исследование коррелированных сигналов в каналах связи с коррелированными сигналами

report submitted for the Confidential Meeting of the Scientific Technological Society of
Radio Engineering and Electrical Communications to A. A. Popov (VSEI), Moscow,
8-10 June. 7-9

AUTHOR: Mityashev, B.N. SOV/109-4-4-10/24
 TITLE: Noise Suppressibility in the Time Discriminators of
 Pulse Signals (O pomekhoustoychivosti vremennogo
 diskriminirovaniya impul'snykh signalov)
 PERIODICAL Radiotekhnika i elektronika, 1959, Vol 4, Nr 4,
 pp 637 - 647 (USSR)

ABSTRACT: In an earlier work (Ref 1), the author carried out an analysis of the pulse signal-plus-noise transfer through a discriminator. The device consisted of two gating stages, two integrators and an adding circuit. The discriminator could be used to determine the position of the signals as a function of the perturbing noise. Here, the accuracy of the instrument is analysed in some detail. When a signal and noise are applied to the discriminator, the slope and the zero of the characteristic of the device are shifted. When the signal-to-noise ratio is small, the average square error of the discriminator is defined by:

$$\Delta \tau_{d0} = \frac{\sigma_d}{S_{d0}} \quad (1)$$

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Noise Suppressibility in the Time Discriminators of Pulse Signals SOV/109-4-4-10/24

where σ_d is the average square value of the noise at the output of the discriminator and S_{d0} is the slope of the characteristic. For large signal-to-noise ratios, a linear approximation can be used and the average square error is given by:

$$\Delta \tau_d = \frac{\sigma_d}{S_d} \quad (2)$$

where $S_d = U_{mcd} / \Delta t_m$ is the average slope of the operating region of the discriminator characteristic. Here, Δt_m is the mismatch of the system, when the output voltage has a maximum value U_{mcd} (Figure 1). When the gating pulses are comparatively long, i.e. when their duration T is longer than the rise time T_a of the signal pulse, the average

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Noise Suppressibility in the Time Discriminators of Pulse Signals

the coefficients b_1 and b_2 are determined by Eqs (10).

The quantity $q = 1/p$ in Eq (9) represents the signal-to-noise ratio at the input of the detector. When the bandwidth of the receiver is comparatively large, the output noise is expressed by Eq (15). The error of the discrimination is therefore given by Eq (16), where K represents the characteristic of the detector, while Δf_0 is the

bandwidth of the receiver. When the duration of the gating pulses is less than the rise time of the signal pulses, the output noise of the discriminator is written as Eq (24)

where σ^2 represents the noise at the input of the system. If the bandwidth is such that $\Delta f_0 T_c$ is greater than 2,

the noise at the output is given by Eq (27). Consequently, for small noise levels, the error in the discrimination is given by Eq (28). In general, the error is expressed by Eq (29). N_0 in Eq (29) denotes the spectral density of

the noise at the input of the filter. The analysis of the

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SOV/109-4-4-10/24
Noise Suppressibility in the Time Discriminators of Pulse Signals

formulae obtained in this work shows that the highest signal-to-noise ratio at the output of the receiving system is obtained when the duration of the gating pulses is equal to that of the signal pulses. When the signal pulses have steep edges, the accuracy of the system can be increased by reducing the duration of the gating pulses. The discrimination accuracy in the case of small noise levels is almost equal to the theoretical maximum, especially when the bandwidth of the filter preceding the detector is inversely proportional to the rise time of the pulse signals and the duration of the gating pulses is equal to the rise time. The author expresses his gratitude to Professor Ye.I. Manayev for constructive criticism of this work. There are 5 figures and 9 references, 1 of which is English and 8 Soviet. One Soviet reference is translated from English.

SUBMITTED: October 16, 1957

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TITLE: Concerning an Optimal Method of P-Code Time-Location

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ABSTRACT: While there are several methods for the determination of time-location of pulsed signals, the maximum accuracy is obtained with the use of the maximum likelihood method. However, this method is complicated and requires a large amount of calculations. The article describes a new method of determining the time-location of pulsed signals in the presence of noise. The method is based on the use of the maximum likelihood method with a simplified algorithm. The results of the calculations are compared with the results of the maximum likelihood method. The results show that the proposed method is simpler and more accurate than the maximum likelihood method. (1) Time Discrimination of P-Code. The results of the calculations are compared with the results of the maximum likelihood method. The results show that the proposed method is simpler and more accurate than the maximum likelihood method. The results show that the proposed method is simpler and more accurate than the maximum likelihood method.

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width T_p of the filter with the optimum value of the parameter A is determined by the condition

$$A = \frac{1}{2} \left(\frac{1}{T_p} + \frac{1}{T_p} \right) = \frac{1}{T_p}$$

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will be considered. The signal is a
A signal, pulse or continuous, is
triangular shape and is located at
the distance r from the origin. It is
that the signal is a pulse of
triangular shape per unit area. The
from the origin is the signal
which takes the shape of a
noise into consideration. The
signal is $\sqrt{\sigma} \sum_{i=1}^N (\sigma \dots)$
detector time). For $\sigma \ll 1$, $K \approx 1$, \approx

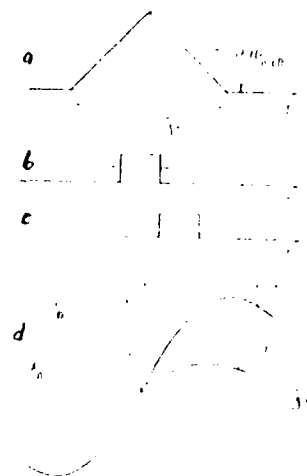
$$\lambda_n(t) = \frac{1}{n}$$

(a) Let the signal be a pulse
of equal duration and area
(b, b). The signal is a pulse of

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Optimal Minimum Mean Squared Error
 Estimation

Difference between the two estimators is
 given by the following expression:

$$\Delta = \frac{1}{n} \sum_{i=1}^n \frac{\partial}{\partial \theta} \log f(\theta; x_i)$$

 where θ is the parameter vector and $f(\theta; x_i)$ is the
 probability density function of the data x_i .
 The limit of the difference between the two
 estimators as $n \rightarrow \infty$ is given by

$$\lim_{n \rightarrow \infty} \Delta = \frac{1}{n} \sum_{i=1}^n \frac{\partial}{\partial \theta} \log f(\theta; x_i)$$

 where θ is the true parameter vector and $f(\theta; x_i)$ is the
 probability density function of the data x_i .
 In this case, the limit of the difference between the two
 estimators is given by the following expression:

$$\lim_{n \rightarrow \infty} \Delta = \frac{1}{n} \sum_{i=1}^n \frac{\partial}{\partial \theta} \log f(\theta; x_i)$$

where σ^2 is the variance of the data x_i and
 the correlation coefficient γ_T and γ_P are given by
 the following expression: $\gamma_T = \frac{1}{n} \sum_{i=1}^n \frac{\partial}{\partial \theta} \log f(\theta; x_i)$
 and $\gamma_P = \frac{1}{n} \sum_{i=1}^n \frac{\partial}{\partial \theta} \log f(\theta; x_i)$.

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$$\tau_T^2 = \frac{2}{T^2} \int_0^T (T - \tau) r(\tau) d\tau, \quad (7)$$

where $r(\tau)$ is correlation coefficient of the signal.
The dependence of coefficients γ_T , γ_{2T} (high and low
ratio noise-to-signal, respectively) and of differences
 $2(\gamma_T'^2 - \gamma_{2T}'^2)$ and $2(\gamma_T''^2 - \gamma_{2T}''^2)$ on T/T_0 are shown in Fig. 2.

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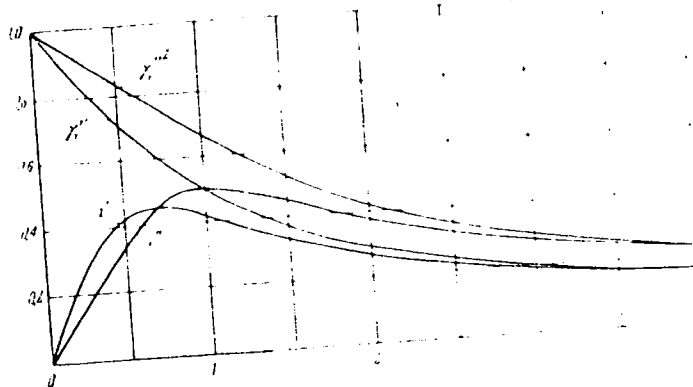


Fig. 2. Coefficients γ_T' , γ_T'' , $a' = (\gamma_T'^2 - \gamma_{2T}'^2)$
and $a'' = 2(\gamma_T''^2 - \gamma_{2T}''^2)$ vs. ratio T/T_0 .

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The effective filter bandwidth (1) is $\Delta f_0 = 1/T_0$ and, therefore, $T/T_0 = \Delta f_0 T$. The RMS time error of tracking caused by the noise is:

$$\Delta \tau_c = \frac{\sigma_d}{S_{d0}} \frac{1}{\sqrt{n}} \quad (13)$$

Here, σ_d is given by Eq. (6), S_{d0} is the discriminator characteristic steepness S_{d0} by Eq. (3); n is number of signal pulses. Since the effective transmission band of filter (1) is $1/T_0$ for the spectral noise density at the filter input N_0 , then $\sigma_d^2 = N_0/T_0$, and therefore

$$\Delta \tau_c = \frac{1}{\sqrt{3}} \sqrt{\frac{N_0 T}{m_{co}}} \frac{1}{\sqrt{n}} \quad (17)$$

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pulses equal to T_p . By a series of calculations, the author obtains:

$$p_{bd} = 0.85 \sqrt{n} \quad (21)$$

$$N_{0\ bd} = 0.37 U_{mc0}^2 T_c \sqrt{n}. \quad (22)$$

Equation (21) is adequate for sufficiently high n , where n is greater than unity. The magnitude of error of low level noises, where $T = T_p$, is:

$$\Delta \tau_c = \frac{1}{2} \sqrt{\frac{N_0 T_c}{U_{mc0}^2}} \sqrt{n} = 0.35 p F \sqrt{n} \quad (23)$$

Thus, a limit accuracy in the discrimination of the signal is different from conditions of a maximum limit of stability. Optimum signal filtration does not provide consolidation of these conditions. But the approximation

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[illegible]
$$\Delta f_0 \approx 1/T_0 \approx 1/T_1$$

filter.
(2) Determination of Pulse Position by Pulse Maximum Value Signals. A determination of the position of pulses by peak values of the signal is made by the following method is proposed: A signal with a maximal filter and detector (Fig. 3a) is limited by a level u_0 and then differentiated (Fig. 3b). The signal is then sent through an amplifier with a low-pass filter. The output pulse has a shape shown in Fig. 3c. The signal is then sent through a differentiating chain (Fig. 3d). The signal is negative pulse separated, or the signal is a positive pulse will be at the output a pulse with a maximum value corresponding with the maximum of the signal pulse (Fig. 3e). The correct setting of a level limiter is important for a high degree of accuracy of low level signal detection.

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high noise stability limit.

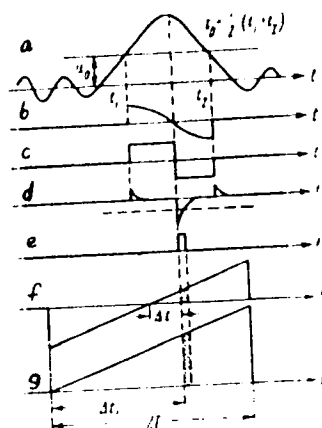


Fig. 5

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The error signal may be developed from a non-linear junction as in Fig. 3e. The error signal, proportional to the mistuning $\Delta\omega$, can be determined by multiplication and a following integration of pulses per Fig. 3f, where the multiplication being easily done if a signal is applied. If it is desirable to register the signal in the form of a voltage, the measuring pulse per Fig. 3f can be used. The authors determine the rms time for targeted signals, pulses and low noise-to-signal ratio per:

$$\Delta\omega = \frac{1}{V} \sqrt{\frac{N_0 T_f}{C_{min}}} \frac{1}{V} \quad (24)$$

The effective pass-band width of an optimal matched filter is within the limits $\Delta\omega_c = (1 - \epsilon) \Delta\omega$, $T_f (\Delta\omega_c = \Delta\omega - \epsilon \Delta\omega)$, $T_f = 0$, and $\Delta\omega_c = \epsilon \Delta\omega$ for $T_f = T_{opt}$. The signal at filter output $U_{out} = (1 - \epsilon) U_{in}$. The rms time can be transformed into:

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